# A Non-Stationary MIMO Vehicle-to-Vehicle Channel Model Based on the Geometrical T-Junction Model

Ali Chelli and Matthias Pätzold

Faculty of Engineering and Science, University of Agder Service Box 509, N-4898 Grimstad, Norway {ali.chelli, matthias.paetzold}@uia.no

*Abstract*—In this paper, we derive a non-stationary multipleinput multiple-output (MIMO) vehicle-to-vehicle (V2V) channel model from the geometrical T-junction model. The propagation environment is assumed to be extremely non-isotropic. The proposed channel model takes into account double-bounce scattering from fixed scatterers. Due to the movement of the transmitter and the receiver, the angles of departure (AoD) and the angles of arrival (AoA) are time-variant, which makes our model nonstationary. In order to study the statistical properties of the proposed channel model, we make use of the Choi-Williams distribution. Analytical solutions are provided for the generalized local autocorrelation function (ACF), the time-frequency distribution, and the local space cross-correlation function (CCF). The proposed channel model is very useful for the design and analysis of future MIMO V2V communication systems.

### I. INTRODUCTION

According to the European commission [1], more than 40 thousand deaths and two million injuries due to road accidents in the European Union have been reported in 2005. It turned out that human errors were involved in 93% of these accidents. V2V communication is a key element in reducing road casualties. For the development of future V2V communication systems, the exact knowledge of the statistics of the underlying fading channel is necessary. Several channel models for V2V communications can be found in the literature. For example, the two-ring channel model for V2V communications has been presented in [2]. There, a reference and a simulation model have been derived starting from the geometrical tworing model. In [3], a three-dimensional reference model for wideband MIMO V2V channels has been proposed. The model takes into account single-bounce and double-bounce scattering in vehicular environments. The geometrical street model [4] captures the propagation effects if the communicating vehicles are moving along a straight street with local roadside obstructions (buildings, trees, etc.). In [5], a statistical frequencyselective channel model for small-scale fading is presented for a V2V communication link.

The majority of channel models that can be found in the literature rely on the stationarity assumption. However, measurement results for V2V channels in [6] have shown that the stationarity assumption is valid only for very short time intervals. This fact arises the need for non-stationary channel models. Actually, if the communicating cars are moving with a relatively high speed, the AoD and the AoA become time-variant resulting in a non-stationary channel model. The traditional framework invoked in case of stationary stochastic processes cannot be used to study the statistical properties of non-stationary channels. In the literature, quite a few time-frequency distributions have been proposed to study non-stationary deterministic signals [7]. A review of these distributions can be found in [7]. Many commonly used timefrequency distributions are members of the Cohen class [8]. It has been stated in [9] that the Cohen class, although introduced for deterministic signals, can be applied on non-stationary stochastic processes.

In this paper, we present a non-stationary MIMO V2V channel model. The AoD and the AoA are supposed to be time dependent. This assumption makes our channel model nonstationary. The correlation properties of a non-stationary channel model can be obtained using a multi-window spectrogram [6]. For rapidly changing spectral content however, finding an appropriate time window size is a rather complicated task. A problem is that a decrease in the time window size improves the time resolution, but reduces the frequency resolution. To overcome this problem, we make use of the Choi-Williams distribution proposed in [10]. The extremely non-isotropic propagation environment is modelled using the T-junction scattering model [11]. In contrast to the original multi-cluster T-model, we assume to simplify matters that each cluster consists of only one scatterer. Under this assumption, the reference and the simulation model are identical. The main contribution of the paper is that it presents a non-stationary channel model with time-variant AoD and AOA. Moreover, analytical expressions for the correlation properties of the nonstationary channel model are provided, evaluated numerically, and then illustrated.

The rest of the paper is organized as follows. In Section II, the geometrical T-model is presented. Based on this geometrical model, we derive a reference (simulation) model in Section III. In Section IV, the correlation properties of the proposed channel model are studied. Numerical results of the correlation functions are presented in Section V. Finally, we draw the conclusions in Section VI.

## II. THE GEOMETRICAL T-JUNCTION MODEL

A typical propagation scenario for V2V communications at a T-junction is presented in Fig. 1. Fixed scatterers are located on both sides of the T-junction. In order to derive the statistical properties of the corresponding MIMO V2V channel, we first need to find a geometrical model that describes properly the vehicular T-junction propagation environment. This geometrical model is illustrated in Fig. 2. It takes into account double-bounce scattering under non-line-of sight conditions. Each building is modelled by one scatterer which makes our model extremely non-isotropic. The scatterers in the neighborhood of the transmitter  $MS_T$  are denoted by  $S_m^T$  $(m = 1, 2, \dots, M)$ , whereas the scatterers close to the receiver  $MS_R$  are designated by  $S_n^R$  (n = 1, 2, ..., N). The total number of scatterers near to the transmitter is denoted by N, while the total number of scatterers near to the receiver is designated by M. The transmitter and the receiver are moving towards the intersection point with the velocities  $v_T$  and  $v_R$ , respectively. The direction of motions of the transmitter and the receiver w.r.t. the x-axis are referred to as  $\phi_T$  and  $\phi_B$ , respectively. The AoD are time-variant and are denoted by  $\alpha_m^T(t)$ , while the symbol  $\beta_n^R(t)$  stands for the AoA. The AoD and the AoA are independent since double-bounce scattering is assumed. The transmitter and the receiver are equipped with an antenna array encompassing  $M_T$  and  $M_R$  antenna elements, respectively. The antenna element spacing at the transmitter side is denoted by  $\delta_T$ . Analogously, the antenna element spacing at the receiver side is referred to as  $\delta_R$ . The tilt angle of the transmit antenna array is denoted by  $\gamma_T$ , while  $\gamma_R$  stands for the tilt angle for the receive antenna array. The transmitter (receiver) is located at a distance  $h_1^T$  ( $h_1^R$ ) from the left-hand side of the street and at a distance  $h_2^T$  ( $h_2^R$ ) from the right-hand side seen in moving direction.



Fig. 1. Typical propagation scenario for V2V communications at a T-junction.

## III. THE REFERENCE MODEL

The starting point for the derivation of the reference model for the MIMO V2V channel is the geometrical T-junction model presented in Fig. 2. For the reference model, we assume double-bounce scattering from fixed scatterers. We distinguish between the scatterers near to the transmitter and the scatterers close to the receiver. It can be seen from Fig. 2 that a wave emitted from the *l*th transmit antenna element  $A_l^T$  ( $l = 1, 2, ..., M_T$ ) travels over the scatterers  $S_m^T$  and  $S_n^R$  before impinging on the *k*th receive antenna element  $A_k^R$ ( $k = 1, 2, ..., M_R$ ). Using the wave propagation model in



Fig. 2. The geometrical T-Junction model for V2V communications.

[2], the complex channel gain  $g_{kl}(\vec{r}_T, \vec{r}_R)$  describing the link  $A_l^T - A_k^R$  of the underlying  $M_T \times M_R$  MIMO V2V channel model can be expressed in the present case as

$$g_{kl}(\vec{r}_T, \vec{r}_R) = \sum_{m=1, n=1}^{M, N} c_{mn} \, e^{j \left(\theta_{mn}(t) + \vec{k}_m^T \cdot \vec{r}_T - \vec{k}_n^R \cdot \vec{r}_R - k_0 d_{mn}(t)\right)}.$$
(1)

The symbols  $c_{mn}$  and  $\theta_{mn}(t)$  stand for the the joint gain and the joint phase shift caused by the scatterers  $S_m^T$  and  $S_n^R$ . The joint channel gain can be written as  $c_{mn} = 1/\sqrt{MN}$  [2]. The phase shift  $\theta_{mn}(t)$  is a stochastic process, as the AoD  $\alpha_m^T(t)$ and the AoA  $\beta_n^R(t)$  are time-variant. This is in contrast to the models proposed in [2] and [11], where the phase shift is a random variable. The joint phase shift can be expressed as  $\theta_{mn}(t) = (\theta_m(t) + \theta'_n(t)) \mod 2\pi$ , where mod stands for the modulo operation. The terms  $\theta_m(t)$  and  $\theta'_n(t)$  are the phase shifts associated with the scatterers  $S_m^T$  and  $S_n^R$ , respectively. The second phase term in (1),  $\vec{k}_m^T \cdot \vec{r}_T$ , is caused by the

The second phase term in (1),  $k_m^T \cdot \vec{r}_T$ , is caused by the movement of the transmitter. The wave vector pointing in the propagation direction of the *m*th transmitted plane wave is denoted by  $\vec{k}_m^T$ , while  $\vec{r}_T$  stands for the spatial translation vector of the transmitter. The scalar product  $\vec{k}_m^T \cdot \vec{r}_T$  can be expanded as

$$\vec{k}_m^T \cdot \vec{r}_T = 2\pi f_{\max}^T \cos(\alpha_m^T(t) - \phi_T)t$$
(2)

where  $f_{\max}^T = \mathbf{v}_T / \lambda$  denotes the maximum Doppler frequency associated with the mobility of the transmitter. The symbol  $\lambda$ refers to the wavelength. The time-variant AoD  $\alpha_m^T(t)$  can be expressed as

$$\alpha_m^T(t) = \begin{cases} -\pi + g_2(t) & \text{if} \quad -\pi \le \alpha_m^T(t) \le -\frac{\pi}{2} \\ g_2(t) & \text{if} \quad -\frac{\pi}{2} < \alpha_m^T(t) \le 0 \\ g_1(t) & \text{if} \quad 0 < \alpha_m^T(t) \le \frac{\pi}{2} \\ \pi + g_1(t) & \text{if} \quad \frac{\pi}{2} < \alpha_m^T(t) \le \pi \end{cases}$$
(3)

where

$$g_{1}(t) = \arctan\left(\frac{h_{1}^{T}\tan(\alpha_{m}^{T}(t_{0}))}{h_{1}^{T}-v_{T}(t-t_{0})\tan(\alpha_{m}^{T}(t_{0}))}\right)$$
(4)

$$g_2(t) = \arctan\left(\frac{h_2 \tan(\alpha_m(t_0))}{h_2^T - \mathbf{v}_T(t - t_0) \tan(\alpha_m^T(t_0))}\right).$$
(5)

We assume that the AoD seen from the transmitter side can be considered as constant for a given time interval if the angular deviation does not exceed a certain threshold. For instance, the AoD  $\alpha_m^T(t_1)$  at time instant  $t_1$  and the AoD  $\alpha_m^T(t_2)$  at time instant  $t_2$  are equal if the angle difference  $|\alpha_m^T(t_1) - \alpha_m^T(t_2)| \le \epsilon_{\alpha}$ , with  $\epsilon_{\alpha}$  is a very small positive value. In this way, the AoD  $\alpha_m^T(t)$  can be written as

$$\alpha_m^T(t) = \alpha_{m,i-1}^T \text{ if } t_{i-1} \le t < t_i \text{ for } i = 1, 2, \dots$$
 (6)

The term  $\alpha_{m,i-1}^T$  is a constant that can be obtained from (3) by setting the time t to  $t_{i-1}$ . The length of the intervals  $[t_{i-1}, t_i)$ and  $[t_i, t_{i+1})$  can be quite different for  $i = 1, 2, \ldots$ . The phase shift introduced by a scatterer is generally dependent on the direction of the outgoing wave. Hence, a change in the AoD  $\alpha_m^T(t)$  results in a new random phase shift. Since the AoD  $\alpha_m^T(t)$  is defined piecewise, the phase shift  $\theta_m(t)$  is also defined piecewise as follows

$$\theta_m(t) = \theta_{m,i-1} \text{ if } t_{i-1} \le t < t_i \text{ for } i = 1, 2, \dots$$
 (7)

where  $\theta_{m,0}$ ,  $\theta_{m,1}$ , ... are independent identically distributed (i.i.d.) random variables uniformly distributed over  $[0, 2\pi)$ .

The third phase term in (1),  $\vec{k}_n^R \cdot \vec{r}_R$ , is associated with the movement of the receiver. The symbol  $\vec{k}_n^R$  stands for the wave vector pointing in the propagation direction of the *n*th received plane wave, while  $\vec{r}_R$  represents the spatial translation vector of the receiver. The scalar product  $\vec{k}_n^R \cdot \vec{r}_R$  can be expanded as

$$\vec{k}_n^R \cdot \vec{r}_R = -2\pi f_{\max}^R \cos(\beta_n^R(t) - \phi_R)t \tag{8}$$

where  $f_{\max}^R = \mathbf{v}_R / \lambda$  denotes the maximum Doppler frequency caused by the receiver movement. Using the geometrical Tjunction model shown in Fig. 2, the time-variant AoA  $\beta_n^R(t)$ can be expressed as

$$\beta_n^R(t) = \begin{cases} -\pi + g_3(t) & \text{if } -\pi \le \beta_n^R(t) \le -\frac{\pi}{2} \\ g_4(t) & \text{if } -\frac{\pi}{2} < \beta_n^R(t) \le \frac{\pi}{2} \\ \pi + g_3(t) & \text{if } \frac{\pi}{2} < \beta_n^R(t) \le \pi \end{cases}$$
(9)

where

$$g_3(t) = \arctan\left(\frac{h_2^R \tan(\beta_n^R(t_0)) - \mathbf{v}_R(t - t_0)}{h_2^R}\right)$$
(10)

$$g_4(t) = \arctan\left(\frac{h_1^R \tan(\beta_n^R(t_0)) - \mathbf{v}_R(t - t_0)}{h_1^R}\right).$$
 (11)

We assume that the AoA seen from the receiver side can be considered as constant for a given time interval if the angular deviation does not exceed a certain threshold. For instance, the AoA  $\beta_n^R(t_1')$  at time instant  $t_1'$  and the AoA  $\beta_n^R(t_2')$  at time instant  $t_2'$  are equal if the angle difference  $|\beta_n^R(t_1') - \beta_n^R(t_2')| \le \epsilon_{\alpha}$ . In this way, the AoA  $\beta_n^R(t)$  can be written as

$$\beta_n^R(t) = \beta_{n,j-1}^R \text{ if } t'_{j-1} \le t < t'_j \text{ for } j = 1, 2, \dots$$
 (12)

The term  $\beta_{n,j-1}^R$  is a constant that can be obtained from (9) by setting the time t to  $t'_{j-1}$ . The length of the intervals  $[t'_{j-1}, t'_j)$ 

and  $[t'_j, t'_{j+1})$  can be quite different for j = 1, 2, ... The phase shift introduced by a scatterer is generally dependent on the direction of the incoming wave. Hence, a change in the AoA  $\beta_n^R(t)$  results in a new random phase shift. Since the AoA  $\beta_n^R(t)$  is defined piecewise, the phase shift  $\theta'_n(t)$  is also defined piecewise as follows

$$\theta'_n(t) = \theta'_{n,j-1} \text{ if } t'_{j-1} \le t < t'_j \text{ for } j = 1, 2, \dots$$
 (13)

where  $\theta'_{n,0}$ ,  $\theta'_{n,1}$ , ... are i.i.d. random variables uniformly distributed over  $[0, 2\pi)$ .

After substituting (2) and (8) in (1), the complex channel gain  $g_{kl}(t)$  can be expressed as

$$g_{kl}(t) = \sum_{m,n=1}^{M,N} \frac{a_m^T b_n^R c_{mn}^{TR}}{\sqrt{MN}} e^{j \left(2\pi (f_m^T + f_n^R)t + \theta_{mn}(t)\right)}$$
(14)

where

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$$a_m^T = e^{j\pi\frac{\delta_T}{\lambda}(M_T - 2l + 1)\cos(\alpha_m^T(t) - \gamma_T)}$$
(15)

$$e_n^R = e^{j\pi\frac{\delta_R}{\lambda}(M_R - 2k + 1)\cos(\beta_n^R(t) - \gamma_R)}$$
(16)

$$\Gamma_{nn}^{R} = e^{-j\frac{2\pi}{\lambda} \left( D_m^T(t) + D_{mn} + D_n^R(t) \right)}$$
(17)

$$f_m^T = f_{\max}^T \cos(\alpha_m^T(t) - \phi^T)$$
(18)

$$f_n^R = f_{\max}^R \cos(\beta_n^R(t) - \phi^R).$$
(19)

with  $D_m^T(t)$  denoting the distance from the transmitter to the scatterer  $S_m^T$ . The term  $D_{mn}$  represents the distance between the scatterers  $S_m^T$  and  $S_n^R$ , while  $D_n^R(t)$  corresponds to the distance from the receiver to the scatterer  $S_n^R$ , as shown in Fig. 2.

### **IV. CORRELATION PROPERTIES**

For wide-sense stationary processes, the temporal ACF depends only on the time difference  $\tau$ . However, for non-stationary processes, the temporal ACF does not only depend on the time difference  $\tau$ , but also on the time t. Due to its time dependance, the ACF of non-stationary processes is called local ACF [7]. Several definitions for the local ACF have been proposed in literature. In this paper, we utilize the definition of the local ACF proposed by Wigner [7], which is given by

$$r_{g_{kl}}(t,\tau) := E\{g_{kl}(t+\tau/2)g_{kl}^*(t-\tau/2)\}$$
(20)

where  $(\cdot)^*$  denotes the complex conjugation and  $E\{\cdot\}$  stands for the expectation operator. By applying the expectation operator on the i.i.d. random variables  $\theta_{m,i}$  (i = 0, 1, ...)and  $\theta'_{n,j}$  (j = 0, 1, ...) and exploiting their independence, we can express the local ACF as

$$r_{g_{kl}}(t,\tau) = r_{g_{kl}}^T(t,\tau) \cdot r_{g_{kl}}^R(t,\tau).$$
(21)

(23)

where

$$r_{g_{kl}}^{T}(t,\tau) = \frac{1}{M} \sum_{m=1}^{M} e^{j2\pi \left(f_{m}^{T}\left(t+\frac{\tau}{2}\right)\left(t+\frac{\tau}{2}\right) - f_{m}^{T}\left(t-\frac{\tau}{2}\right)\left(t-\frac{\tau}{2}\right)\right)}$$

$$r_{g_{kl}}^{R}(t,\tau) = \frac{1}{N} \sum_{n=1}^{N} e^{j2\pi \left(f_{n}^{R}\left(t+\frac{\tau}{2}\right)\left(t+\frac{\tau}{2}\right) - f_{n}^{R}\left(t-\frac{\tau}{2}\right)\left(t-\frac{\tau}{2}\right)\right)}.$$
(22)

Note that the local ACF  $r_{g_{kl}}(t,\tau)$  is written as a product of the local transmit ACF  $r_{g_{kl}}^T(t,\tau)$  and the local receive ACF  $r_{q_{kl}}^T(t,\tau)$  since we assume a limited number of scatterers in the proposed model. The expression of the local ACF is derived using the Wigner method [7]. By applying the Fourier transformation on the local ACF in (21), we obtain the Wigner time-frequency distribution. The former even though a member of the Cohen class of distributions, suffers from the cross-term problem [7]. To deal with this issue, a kernel function aiming to reduce the cross-terms need to be introduced. One of the effective distributions in diminishing the effect of cross-terms is the Choi-Williams distribution [10]. Choi and Williams devised their kernel function in such a way that a relatively large weight is given to  $g_{kl}(u+\tau/2)g_{kl}^*(u-\tau/2)$  if u is close to t. In this way, they emphasis the local behaviour of the channel and guarantee that the non-stationarities will not be smeared in time and frequency. The kernel function for the Choi-Williams distribution is given by  $\phi(\xi, \tau) = e^{-\xi^2 \tau^2/\sigma}$ . It follows that the generalized local ACF can be expressed as [10]

$$\begin{split} K(t,\tau;\phi) &= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{j2\pi\xi(u-t)} \phi(\xi,\tau) r_{g_{kl}}(u,\tau) \, du \, d\xi \\ &= \int_{-\infty}^{\infty} \frac{r_{g_{kl}}(u,\tau)}{\sqrt{4\pi\tau^2/\sigma}} \exp\left(-\frac{(u-t)^2}{4\tau^2/\sigma}\right) \, du. \ (24) \end{split}$$

The generalized local ACF presented above can be used for both stationary and non-stationary processes. Actually, if the process is stationary, the local ACF  $r_{g_{kl}}(u,\tau)$  equals  $r_{g_{kl}}(\tau)$ . Using (24), it turns out that the generalized local ACF for stationary processes equals the classical ACF, i.e.,  $K(t,\tau;\phi) = r_{g_{kl}}(\tau)$ .

For stationary processes, the power spectral density can be obtained from the Fourier transformation of the temporal ACF. Analogously, for non-stationary processes, the time-frequency distribution can be obtained from the generalized local ACF by applying the Fourier transformation. The time-frequency distribution gives an insight into how the power spectrum varies with time t. The time-frequency distribution  $W(t, f; \phi)$  can be written as

$$W(t,f;\phi) = \int_{-\infty}^{\infty} K(t,\tau;\phi) e^{-j2\pi f\tau} d\tau$$
$$= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} A(\xi,\tau;\phi) e^{-j2\pi (f\tau - t\xi)} d\tau d\xi$$
(25)

where  $A(\xi, \tau; \phi)$  is the ambiguity function.

The local space CCF can be expressed as

$$\rho_{kl,k'l'}(t,\delta_T,\delta_R) = E\{g_{kl}(t)g_{k'l'}^*(t)\}$$

$$= \frac{1}{M} \sum_{m=1}^M e^{-j2\pi \frac{\delta_T}{\lambda}(l-l')\cos(\alpha_m^T(t)-\gamma_T)}$$

$$\cdot \frac{1}{N} \sum_{n=1}^N e^{-j2\pi \frac{\delta_R}{\lambda}(k-k')\cos(\beta_n^R(t)-\gamma_R)}$$

$$= \rho_{ll'}^{T}(t,\delta_T) \cdot \rho_{kk'}^R(t,\delta_R).$$
(26)

In (26), the AoD  $\alpha_m^T(t)$  and the AoA  $\beta_n^R(t)$  are given by (6) and (12), respectively. Note that the local space CCF is written as a product of the local transmit space correlation function (CF)  $\rho_{ll'}^T(t, \delta_T)$  and the local receive space CF  $\rho_{kk'}^R(t, \delta_R)$ .

## V. NUMERICAL RESULTS

In this section, the analytical expressions presented in the previous section are evaluated numerically and then illustrated. The propagation environment encompasses twelve scatterers around the transmitter. Six scatterers are located on the left side of the transmitter and the remaining scatterers are on the right side. The distance between two successive scatterers is set to 20 m. We consider the same number of scatterers around the receiver. The transmitter and the receiver have a velocity of 70 km/h and a direction of motion determined by  $\phi_T = 0$  and  $\phi_R = -\pi/2$ , respectively. The transmitter and the receiver antenna tilt angles  $\gamma_T$  and  $\gamma_R$  are equal to  $\pi/2$ . The street parameters are chosen as  $h_1^T = h_2^T = 50$  m and  $h_1^R = h_2^R = 50$  m. The parameter  $\epsilon_{\alpha}$  is set to 0.1.

The absolute value of the resulting generalized local ACF  $K(t, \tau; \phi)$  is illustrated in Fig. 3. From this figure, we can see that the shape of the local ACF changes for different values of t, which is due to the non-stationarity of the channel model. If the channel model is stationary, we would observe the same shape of the local ACF at different time instants t. The absolute value of the time-frequency distribution,  $|W(t, f; \phi)|$ , shown in Fig. 4, is obtained from the generalized local ACF by applying the Fourier transform w.r.t. the time lag  $\tau$ . It can be seen from this figure how the Doppler spectrum of the channel varies with time t. For the chosen propagation scenario, it can be observed from Fig. 4 that the zero Doppler frequency has the highest power for all time instants t. The power of the non-zero Doppler frequencies decays for a certain period of time before increasing again. The absolute value of the local transmit space CF  $|\rho_{12}^T(t, \delta_T)|$  is presented in Fig. 5. It can be seen from this figure that the amplitude of this function is more sensitive to the transmit antenna spacing  $\delta_T$  than to the time t. The absolute value of the local receive space CF  $|\rho_{12}^R(t,\delta_R)|$  is illustrated in Fig. 6. For the chosen scenario, this function decays faster than the local transmit space CF w.r.t. the antenna spacing.



Fig. 3. The absolute value of the generalized local ACF  $|K(t, \tau; \phi)|$ .



Fig. 4. The absolute value of the time-frequency distribution  $|W(t, f; \phi)|$ .



Fig. 5. Absolute value of the local transmit space CF  $|\rho_{12}^T(t, \delta_T)|$ .

## VI. CONCLUSION

In this paper, we have presented a non-stationary MIMO V2V channel model. Based on the geometrical T-junction model, we have derived an expression for the time-variant channel gain taking into account double-bounce scattering from fixed scatterers. We have assumed a limited number of scatterers. Under this assumption, the reference model equals the simulation model. In vehicular environments, the high speed of the communicating vehicles results in time-variant AoD and AoA. This property is taken into account in our channel model, which makes the model non-stationary. To study the statistical properties of the proposed channel model, we utilized the Choi-Williams distribution. We have provided analytical expressions of the generalized local ACF, the timefrequency distribution, and the local space CCF. The latter can be written as a product of the local transmit space CF and the local receive space CF. Supported by our analysis, we can conclude that the stationarity assumption is violated for V2V channels, especially if the mobile speed is high and the observation interval is large. Non-stationary channel models are needed as a tool for designing future V2V communication systems. In future work, the effect of moving scatterers on the channel statistics will be studied.



Fig. 6. Absolute value of the local receive space CF  $|\rho_{12}^R(t, \delta_R)|$ .

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